

**Warsaw University
of Technology**



**Faculty of Power and
Aeronautical Engineering**

WARSAW UNIVERSITY OF TECHNOLOGY

Institute of Aeronautics and Applied Mechanics

Finite element method 2 (FEM 2)

Introduction to structural dynamics

10.2021

INTRODUCTION TO STRUCTURAL DYNAMICS

The transient dynamic equilibrium equation for a discrete FE model:

$$\begin{array}{ccccccc} [M] \cdot \{\ddot{q}\} & + & [C] \cdot \{\dot{q}\} & + & [K] \cdot \{q(t)\} & = & \{F(t)\} \\ \text{NDOF} \times \text{NDOF} & & \text{NDOF} \times 1 & & \text{NDOF} \times \text{NDOF} & & \text{NDOF} \times 1 \\ & & \uparrow & & \uparrow & & \uparrow \\ & & \text{VECTOR} & & \text{VECTOR} & & \text{VECTOR} \\ & & \text{OF NODAL} & & \text{OF NODAL} & & \text{OF NODAL} \\ & & \text{ACCELERATIONS} & & \text{VELOCITIES} & & \text{DISPLACEMENTS} \\ & & & & & & \uparrow \\ \text{MASS} & & \text{DAMPING} & & \text{STIFFNESS} & & \text{LOAD} \\ \text{MATRIX} & & \text{MATRIX} & & \text{MATRIX} & & \text{VECTOR} \end{array}$$

$$\underbrace{\{\dot{q}\}}_{\text{NDOF} \times 1} = \frac{d}{dt} \underbrace{\{q(t)\}}_{\text{NDOF} \times 1} = \left\{ \begin{array}{c} \frac{dq_1}{dt} \\ \frac{dq_2}{dt} \\ \vdots \\ \frac{dq_{\text{NDOF}}}{dt} \end{array} \right\}$$

$$\underbrace{\{\ddot{q}\}}_{\text{NDOF} \times 1} = \frac{d}{dt} \underbrace{\{\dot{q}\}}_{\text{NDOF} \times 1} = \left\{ \begin{array}{c} \frac{d^2q_1}{dt^2} \\ \frac{d^2q_2}{dt^2} \\ \vdots \\ \frac{d^2q_{\text{NDOF}}}{dt^2} \end{array} \right\}$$

SPECIAL CASES :

* HARMONIC VIBRATIONS :

$$\underbrace{\{F(t)\}}_{\text{NDOF} \times 1} = \left\{ \begin{array}{l} F_{a_1} \cdot \sin(\omega t + \varphi_1) \\ F_{a_2} \cdot \sin(\omega t + \varphi_2) \\ \vdots \\ F_{a_{\text{NDOF}}} \cdot \sin(\omega t + \varphi_{\text{NDOF}}) \end{array} \right\}$$

↑
load
magnitudes

↑
phase
shifts

ω - angular frequency , $\omega = 2\pi \cdot f$ $[\frac{1}{s}]$

f - frequency $[\frac{1}{s}]$

* FREE VIBRATIONS :

$$\underbrace{\{F(t)\}}_{\text{NDOF} \times 1} = \underbrace{\{0\}}_{\text{NDOF} \times 1}$$

* FREE UNDAMPED VIBRATIONS (NATURAL VIBRATIONS) :

$$\underbrace{\{F(t)\}}_{\text{NDOF} \times 1} = \underbrace{\{0\}}_{\text{NDOF} \times 1} \text{ AND } \underbrace{[C]}_{\text{NDOF} \times \text{NDOF}} = \underbrace{[0]}_{\text{NDOF} \times \text{NDOF}}$$

* STATIC ANALYSIS :

$$\underbrace{[M]}_{\text{NDOF} \times \text{NDOF}} = \underbrace{[0]}_{\text{NDOF} \times \text{NDOF}}, \quad \underbrace{[C]}_{\text{NDOF} \times \text{NDOF}} = \underbrace{[0]}_{\text{NDOF} \times \text{NDOF}}$$

$$\underbrace{\{F(t)\}}_{\text{NDOF} \times 1} = \underbrace{\{F\}}_{\text{NDOF} \times 1}; \quad F_1, F_2, \dots = \text{const}$$

NATURAL VIBRATIONS

$$\begin{matrix} [M] & \cdot & \{\ddot{q}\} & + & [K] & \cdot & \{q(t)\} & = & \{0\} \\ \text{NDOF} \times \text{NDOF} & & \text{NDOF} \times 1 & & \text{NDOF} \times \text{NDOF} & & \text{NDOF} \times 1 & & \text{NDOF} \times 1 \end{matrix}$$

general solution :

$$\begin{matrix} \{q(t)\} & = & \{q\}_A & \cdot \cos \omega t & + & \{q\}_B & \cdot \sin \omega t \\ \text{NDOF} \times 1 & & \text{NDOF} \times 1 & & & \text{NDOF} \times 1 & \\ & & \uparrow & & & \swarrow & \\ & & \text{depend on initial conditions} & & & & \end{matrix}$$

VELOCITY VECTOR :

$$\underbrace{\{ \dot{q}(t) \}}_{\text{NDOF} \times 1} = -\omega \underbrace{\{ q \}}_A \cdot \sin \omega t + \omega \underbrace{\{ q \}}_B \cdot \cos \omega t$$

NDOF × 1 NDOF × 1 NDOF × 1

ACCELERATION VECTOR :

$$\underbrace{\{ \ddot{q}(t) \}}_{\text{NDOF} \times 1} = -\omega^2 \underbrace{\{ q \}}_A \cdot \cos \omega t - \omega^2 \underbrace{\{ q \}}_B \sin \omega t =$$
$$= -\omega^2 \underbrace{\{ q(t) \}}_{\text{NDOF} \times 1}$$

$$\begin{array}{ccccccc}
 [M] \cdot (-\omega^2 \{q(t)\}) + [K] \cdot \{q(t)\} = \{0\} \\
 \text{NDOF} \times \text{NDOF} & \text{NDOF} \times 1 & \text{NDOF} \times \text{NDOF} & \text{NDOF} \times 1 & \text{NDOF} \times 1
 \end{array}$$

eigenvalue problem

$$\left(\begin{array}{c} [K] \\ \text{NDOF} \times \text{NDOF} \end{array} - \omega^2 \begin{array}{c} [M] \\ \text{NDOF} \times \text{NDOF} \end{array} \right) \{q(t)\} = \{0\}$$

nontrivial solution :

- an eigenvector $\{q(t)\} \neq \{0\}$
- $\det \left(\begin{array}{c} [K] \\ \text{NDOF} \times \text{NDOF} \end{array} - \omega^2 \begin{array}{c} [M] \\ \text{NDOF} \times \text{NDOF} \end{array} \right) = 0$

COMMENT 1 :

if $\det([K] - \omega^2[M]) \neq 0$ then

the inverse matrix $([K] - \omega^2[M])^{-1}$ exists :

$$\underbrace{([K] - \omega^2[M])^{-1} \cdot ([K] - \omega^2[M])}_{= [I] - \text{identity matrix}} \cdot \{q(t)\} = \underbrace{([K] - \omega^2[M])^{-1} \cdot \{0\}}_{= \{0\}}$$

thus :

$$[I] \cdot \{q(t)\} = \{q(t)\} = \{0\}$$

what is in contradiction to $\{q(t)\} \neq \{0\}$

COMMENT 2 : IF $\det [M] \neq 0 \Rightarrow$

$$[M]^{-1} \cdot ([K] - \omega^2 [M]) \cdot \{q(t)\} = [M]^{-1} \cdot \{0\}$$

$$\underbrace{([M]^{-1} [K] - \omega^2 [I])}_{[A]} \{q(t)\} = \{0\}$$

$$[A] , \quad \omega^2 = c , \quad \{q(t)\} = \{v\}$$

$$[A] \cdot \{v\} = c \cdot \{v\}$$



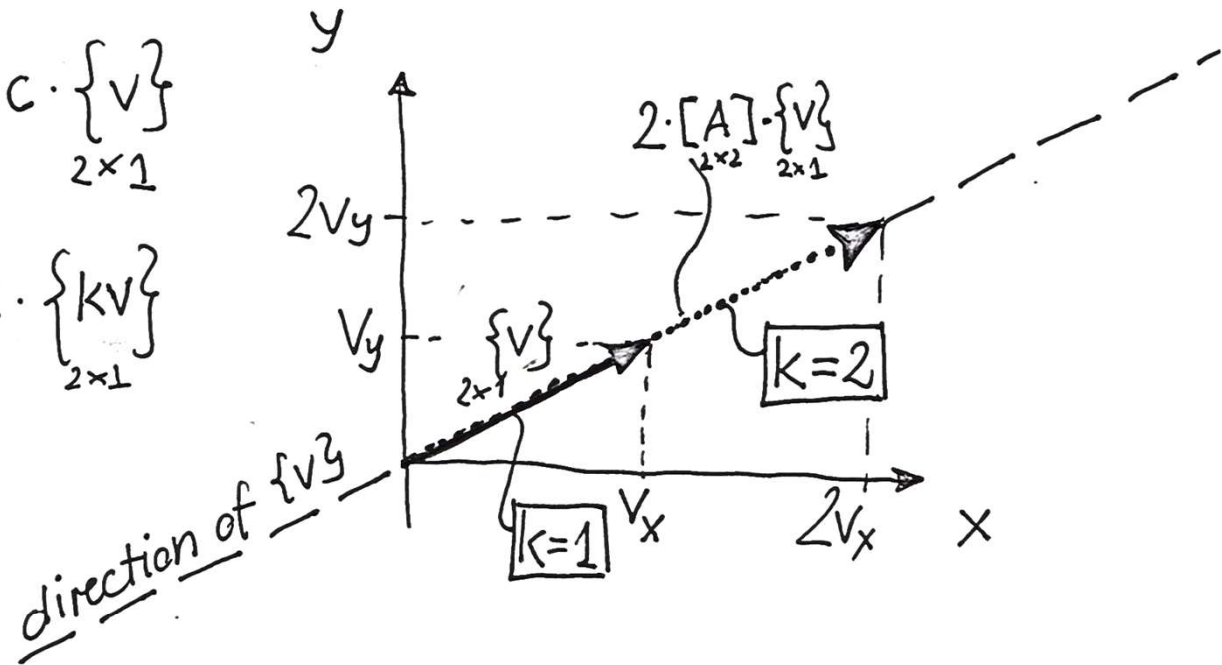
eigenvalue of $[A]$

eigenvector of $[A]$ corresponding to the eigenvalue c

FOR ANY CONSTANT k ($NDOF = 2$) :

$$k \cdot \underset{2 \times 2}{[A]} \cdot \underset{2 \times 1}{\{v\}} = k \cdot c \cdot \underset{2 \times 1}{\{v\}}$$

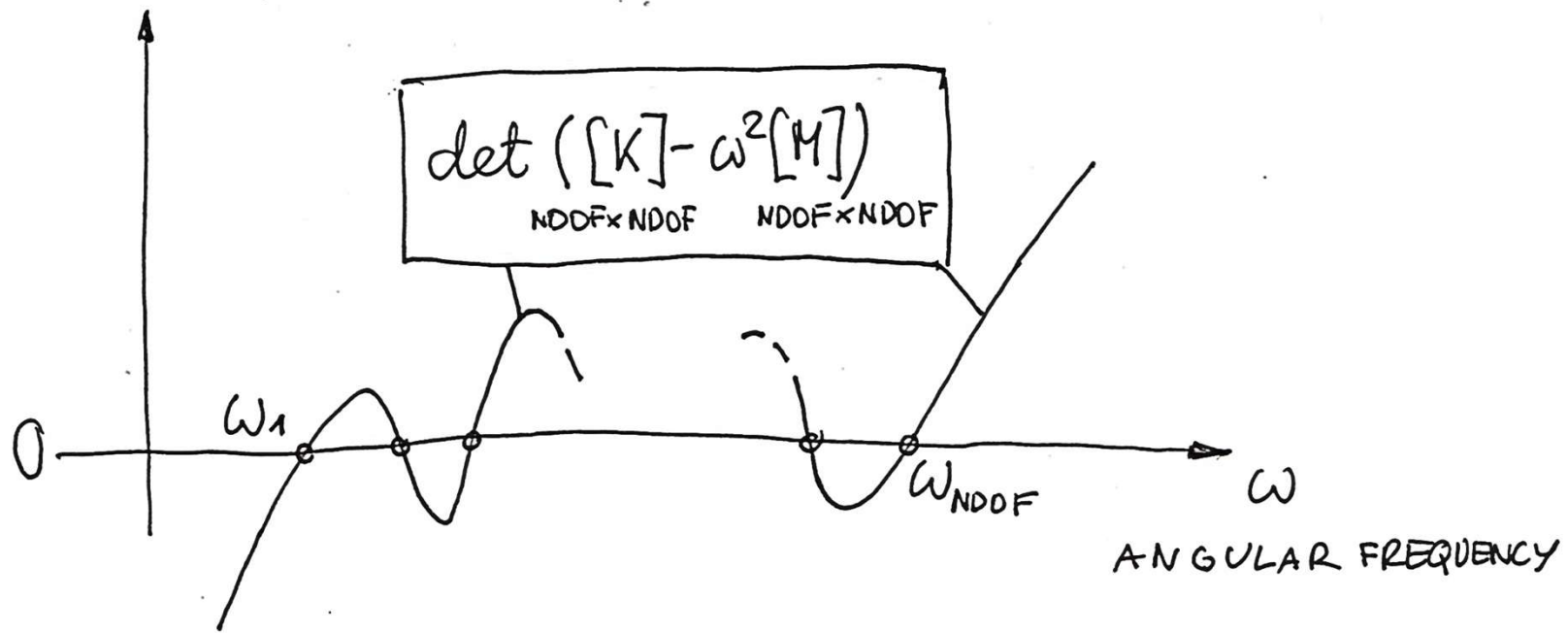
$$\underset{2 \times 2}{[A]} \cdot \underset{2 \times 1}{\{kv\}} = c \cdot \underset{2 \times 1}{\{kv\}}$$



Conclusions :

1) $\frac{V_y}{V_x} = \text{const} \Rightarrow V_x$ and V_y are dependent on each other

2) $\{kv\}$ is also the solution of an eigenvalue problem



$\omega_1, \omega_2, \dots, \omega_i, \dots, \omega_{NDOF}$ - NATURAL ANGULAR FREQUENCIES

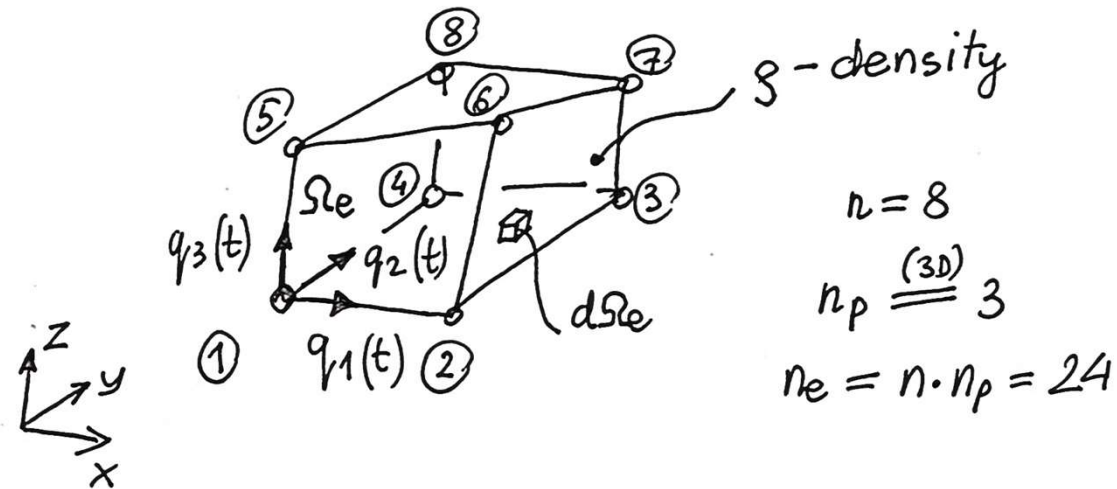
The shape of vibration (the vibration mode) for ω_i is described by the corresponding eigenvector $\{q\}_i$, which is normalized :

$$\underbrace{[q]_i}_{1 \times \text{NDOF}} \cdot \underbrace{[I]}_{\text{NDOF} \times \text{NDOF}} \cdot \underbrace{\{q\}_i}_{\text{NDOF} \times 1} = 1 \quad \text{or} \quad \underbrace{[q]_i}_{1 \times \text{NDOF}} \cdot \underbrace{[M]}_{\text{NDOF} \times \text{NDOF}} \cdot \underbrace{\{q\}_i}_{\text{NDOF} \times 1} = 1$$

ANALYSIS OF NATURAL VIBRATIONS (MODAL ANALYSIS)

- is more time-consuming than a linear structural static analysis,
- iterative solver is used for the limited number of ω_i ,
- the effect of stress can be included (prestress).

MASS MATRIX OF A FINITE ELEMENT



LOCAL VECTOR OF NODAL PARAMETERS :

$$\{q(t)\}_e = [q_1(t), q_2(t), \dots, q_{n_e}(t)]_e$$

$1 \times n_e$

DISPLACEMENT VECTOR :

$$\begin{Bmatrix} u(t) \\ v(t) \\ w(t) \end{Bmatrix}_{3 \times 1} = [N]_{3 \times n_e} \cdot \begin{Bmatrix} q(t) \end{Bmatrix}_e_{n_e \times 1}$$

VELOCITY VECTOR:

$$\begin{matrix} \{\dot{u}\} \\ 3 \times 1 \end{matrix} = \begin{matrix} [N] \\ 3 \times n_e \end{matrix} \cdot \begin{matrix} \{\dot{q}\}_e \\ n_e \times 1 \end{matrix} \quad , \quad \begin{matrix} [\dot{u}] \\ 1 \times 3 \end{matrix} = \begin{matrix} [\dot{q}]_e \\ 1 \times n_e \end{matrix} \cdot \begin{matrix} [N]^T \\ n_e \times 3 \end{matrix}$$

KINETIC ENERGY OF AN INFINITESIMALLY SMALL PART $d\Omega_e$:

$$dT_e = \frac{1}{2} \begin{matrix} [\dot{u}] \\ 1 \times 3 \end{matrix} dm \cdot \begin{matrix} \{\dot{u}\} \\ 3 \times 1 \end{matrix} = \frac{1}{2} \begin{matrix} [\dot{q}]_e \\ 1 \times n_e \end{matrix} \begin{matrix} [N]^T \\ n_e \times 3 \end{matrix} \rho d\Omega_e \begin{matrix} [N] \\ 3 \times n_e \end{matrix} \begin{matrix} \{\dot{q}\}_e \\ n_e \times 1 \end{matrix} =$$

$$= \frac{1}{2} \begin{matrix} [\dot{q}]_e \\ 1 \times n_e \end{matrix} \cdot \rho \begin{matrix} [N]^T \\ n_e \times 3 \end{matrix} \cdot \begin{matrix} [N] \\ 3 \times n_e \end{matrix} \cdot \begin{matrix} \{\dot{q}\}_e \\ n_e \times 1 \end{matrix} d\Omega_e$$

KINETIC ENERGY OF A FINITE ELEMENT:

$$T_e = \int dT_e = \int_{\Omega_e} \frac{1}{2} \underbrace{L \dot{q}}_{1 \times n_e} \underbrace{\rho}_{\text{scalar}} \underbrace{[N]^T}_{n_e \times 3} \underbrace{[N]}_{3 \times n_e} \cdot \underbrace{\{\dot{q}\}_e}_{n_e \times 1} d\Omega_e =$$

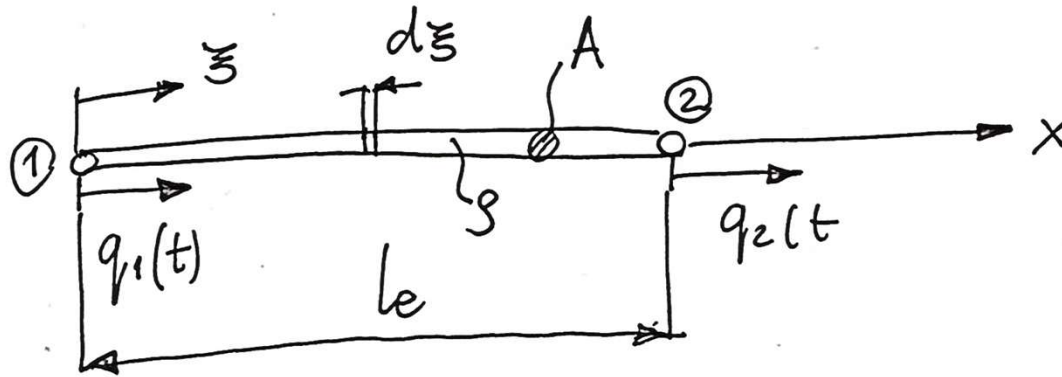
$$= \frac{1}{2} \underbrace{L \dot{q}}_{1 \times n_e} \cdot \underbrace{\int_{\Omega_e} \rho [N]^T [N] d\Omega_e}_{n_e \times n_e} \cdot \underbrace{\{\dot{q}\}_e}_{n_e \times 1} =$$

$[m]_e$ — consistent mass matrix
 $n_e \times n_e$

$$= \frac{1}{2} \underbrace{L \dot{q}}_{1 \times n_e} \cdot \underbrace{[m]_e}_{n_e \times n_e} \cdot \underbrace{\{\dot{q}\}_e}_{n_e \times 1}$$

A lumped mass matrix $[m_L]_e = \frac{\rho \Omega_e}{n_e} \cdot [I]_{n_e \times n_e}$

MASS MATRIX OF A BAR ELEMENT



$$d\Omega_e = A \cdot d\xi, \quad dm = \rho d\Omega_e = \underbrace{\rho A}_{\text{linear density}} \cdot d\xi \quad \left[\frac{\text{kg}}{\text{m}} \right]$$

$$u(\xi, t) = [N_1(\xi), N_2(\xi)] \cdot \begin{Bmatrix} q_1(t) \\ q_2(t) \end{Bmatrix}_e$$

$$N_1(\xi) = 1 - \frac{\xi}{l_e}, \quad N_2(\xi) = \frac{\xi}{l_e}$$

$$\dot{u}(\xi) = [N_1, N_2] \cdot \begin{Bmatrix} \dot{q}_1 \\ \dot{q}_2 \end{Bmatrix}_e = [\dot{q}_1, \dot{q}_2] \cdot \begin{Bmatrix} N_1 \\ N_2 \end{Bmatrix}$$

$$dT_e = \frac{1}{2} \dot{u}(\xi) \cdot dm \cdot \dot{u}(\xi) = \frac{1}{2} \underbrace{[\dot{q}]_e}_{1 \times 2} \cdot \begin{Bmatrix} N_1 \\ N_2 \end{Bmatrix} \rho A \cdot [N_1, N_2] \cdot \underbrace{\{\dot{q}\}_e}_{2 \times 1} d\xi$$

$$T_e = \int dT_e = \frac{1}{2} \underbrace{[\dot{q}]_e}_{1 \times 2} \cdot \rho A \cdot \underbrace{\int_0^l \begin{bmatrix} N_1 N_1 & N_1 N_2 \\ N_2 N_1 & N_2 N_2 \end{bmatrix} d\xi}_{[M]_e} \cdot \underbrace{\{\dot{q}\}_e}_{2 \times 1}$$

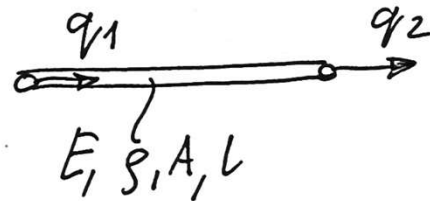
$$\int_0^l N_1 N_1 d\xi = \int_0^l \left(1 - \frac{\xi}{l}\right)^2 d\xi = \int_0^l \left(1 - \frac{2\xi}{l} + \frac{\xi^2}{l^2}\right) d\xi = l - \frac{l^2}{l} + \frac{l^3}{3l^2} = \frac{l}{3}$$

$$[m]_e = \frac{\rho A l}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix}, \quad [m_L]_e = \frac{\rho A l}{2} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

EXAMPLE. FIND NATURAL ANGULAR FREQUENCIES AND VIBRATION MODES FOR A SINGLE BAR ELEMENT.

CONSIDER UNCONSTRAINED AND CONSTRAINED BAR.

UNCONSTRAINED BAR:



$$[k]_e = \frac{EA}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} = [K]_{2 \times 2}$$

$$[m]_e = \frac{\rho AL}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} = [M]_{2 \times 2}$$

$$\left(\begin{matrix} [K]_{2 \times 2} \\ -\omega^2 [M]_{2 \times 2} \end{matrix} \right) \{q\}_{2 \times 1} = \{0\}_{2 \times 1}$$

$$\det \left(\begin{matrix} [K]_{2 \times 2} \\ -\omega^2 [M]_{2 \times 2} \end{matrix} \right) = 0$$

$$\det \left(\frac{EA}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} - \omega^2 \cdot \frac{\rho AL}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \right) = 0$$

$$\frac{EA}{L} \cdot \det \left(\begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} - \frac{\omega^2 \rho L^2}{6E} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \right) = 0$$

$$\lambda = \frac{\omega^2 \rho L^2}{6E} \Rightarrow \det \begin{bmatrix} 1-2\lambda & -(1+\lambda) \\ -(1+\lambda) & 1-2\lambda \end{bmatrix} = 0$$

$$(1-2\lambda)^2 - (1+\lambda)^2 = 0$$

$$1 - 4\lambda + 4\lambda^2 - 1 - 2\lambda - \lambda^2 = 0$$

$$3\lambda^2 - 6\lambda = 0$$

$$3\lambda(\lambda - 2) = 0$$

$$\Downarrow \quad \Downarrow$$

$$\lambda_1 = 0 \quad \lambda_2 = 2$$

$$\omega_i = \sqrt{\frac{6E\lambda_i}{\rho l^2}}$$

$$\omega_1 = \sqrt{\frac{6E\lambda_1}{\rho l^2}} = 0$$

$$\omega_2 = \sqrt{\frac{6E\lambda_2}{\rho l^2}} = \frac{2\sqrt{3}}{l} \sqrt{\frac{E}{\rho}}$$

VIBRATION MODES

$$\frac{EA}{l} \left(\begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} - \lambda \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \right) \cdot \{q\} = \{0\}$$

$$\begin{bmatrix} 1-2\lambda & -(1+\lambda) \\ -(1+\lambda) & 1-2\lambda \end{bmatrix} \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

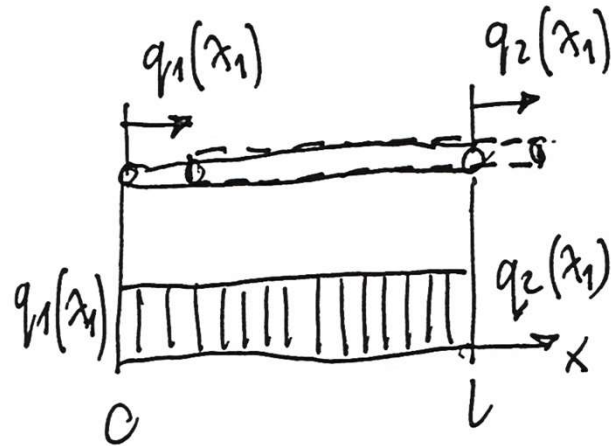
$$\lambda_1 = 0 :$$

$$\begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} q_1(x_1) \\ q_2(x_1) \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

$$\left. \begin{array}{l} 1 \cdot q_1(x_1) - 1 \cdot q_2(x_1) = 0 \\ -1 \cdot q_1(x_1) + 1 \cdot q_2(x_1) = 0 \end{array} \right\} \text{equations are linearly dependent}$$

$$q_1(x_1) = q_2(x_1)$$

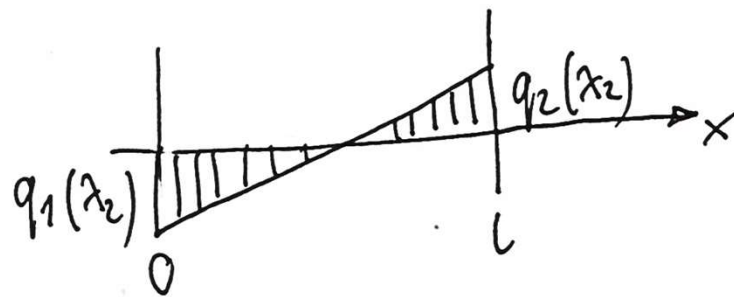
MOVEMENT AS A RIGID BODY



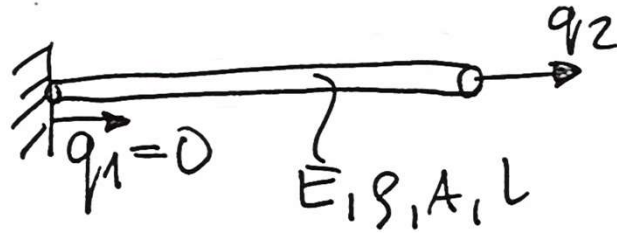
$$\lambda_2 = 2 :$$

$$\begin{bmatrix} -3 & -3 \\ -3 & -3 \end{bmatrix} \cdot \begin{Bmatrix} q_1(\lambda_2) \\ q_2(\lambda_2) \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

$$-3q_1(\lambda_2) - 3q_2(\lambda_2) = 0 \Rightarrow q_2(\lambda_2) = -q_1(\lambda_2)$$



CONSTRAINED BAR :



$$\left(\begin{bmatrix} H & \\ & \text{hatched} \end{bmatrix} - \omega^2 \begin{bmatrix} H & \\ & \text{hatched} \end{bmatrix} \right) \begin{Bmatrix} q_1 \\ q_2 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

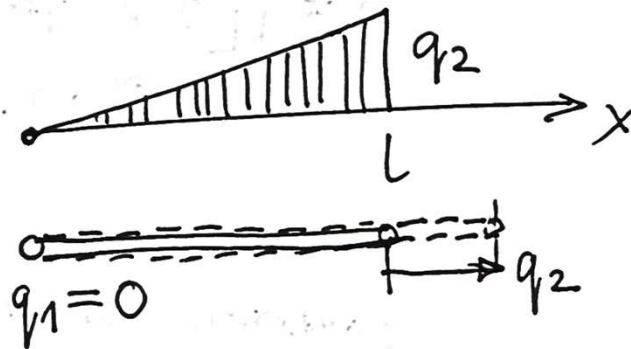
$$\left(\frac{EA}{L} - \omega_1^2 \frac{\rho AL}{3} \right) \cdot q_2 = 0$$

$$\omega_1 = \frac{\sqrt{3}}{L} \sqrt{\frac{E}{\rho}}$$

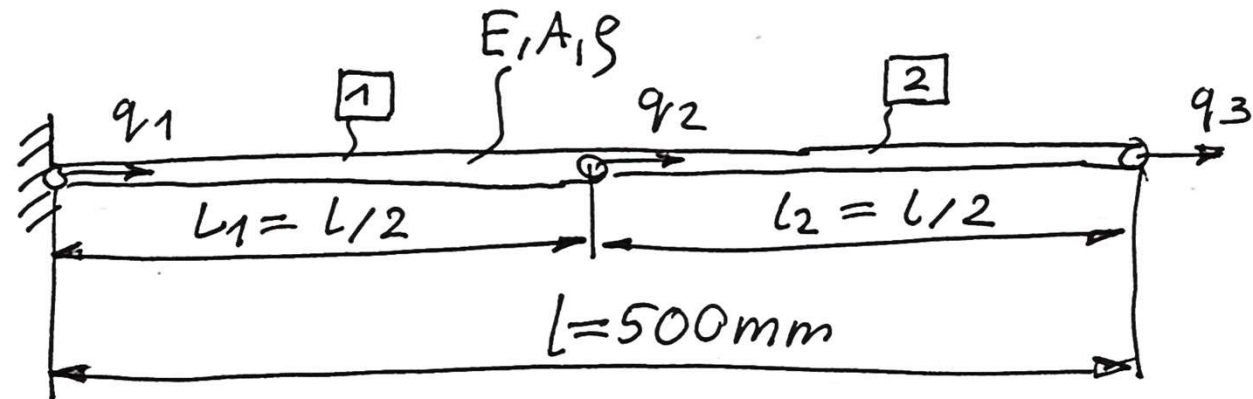
Analytical solution : $\bar{\omega}_i = \frac{(i-0.5)\pi}{L} \sqrt{\frac{E'}{S}}$

$$\bar{\omega}_1 = \frac{1.5708}{L} \sqrt{\frac{E'}{S}}$$

Relative error: $\Delta\omega_1 = \frac{\omega_1 - \bar{\omega}_1}{\bar{\omega}_1} = 10\%$



EXAMPLE. FIND NATURAL FREQUENCIES AND VIBRATION MODES FOR A BAR CONSTRAINED AT ONE END USE 2 FES.



$$E = 2 \cdot 10^5 \text{ MPa}$$

$$\rho = 7.8 \cdot 10^3 \frac{\text{kg}}{\text{m}^3} = 7.8 \cdot 10^3 \frac{\text{kg m}}{\text{s}^2 \text{m}^3} \cdot \frac{\text{s}^2}{\text{m}} = \left. \begin{array}{l} \\ \\ \end{array} \right\} \text{STEEL}$$

$$= 7.8 \cdot 10^3 \frac{\text{Ns}^2}{\text{m}^4} = 7.8 \cdot 10^{-9} \frac{\text{Ns}^2}{\text{mm}^4}$$

$$A = 100 \text{ mm}^2, \quad l_1 = l_2 = l/2$$

FE solution :

$$[q]_{1 \times 3} = [q_1, q_2, q_3]$$

$$[K]_1 = \frac{EA}{L_1} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} = [K]_2 = \frac{2EA}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix}$$

$$[m]_1 = [m]_2 = \frac{\rho A L e}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} = \frac{\rho A L}{12} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix}$$

$$[K]_{3 \times 3} = \begin{bmatrix} [K]_1 & 0 \\ 0 & [K]_2 \end{bmatrix} = \frac{2EA}{L} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 1 \end{bmatrix}$$

$$[M]_{3 \times 3} = \begin{bmatrix} [m]_1 & 0 \\ 0 & [m]_2 \end{bmatrix} = \frac{\rho A L}{12} \begin{bmatrix} 2 & 1 & 0 \\ 1 & 4 & 1 \\ 0 & 1 & 2 \end{bmatrix}$$

$$\left(\begin{matrix} [K] & -\omega^2 [M] \\ 3 \times 3 & 3 \times 3 \end{matrix} \right) \cdot \begin{matrix} \{q\} \\ 3 \times 1 \end{matrix} = \begin{matrix} \{0\} \\ 3 \times 1 \end{matrix}, \quad q_1 = 0$$

$$\left(\frac{2EA}{L} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 1 \end{bmatrix} - \omega^2 \cdot \frac{\rho AL}{12} \begin{bmatrix} 2 & 1 & 0 \\ 1 & 4 & 1 \\ 0 & 1 & 2 \end{bmatrix} \right) \cdot \begin{Bmatrix} 0 \\ q_2 \\ q_3 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix}$$

$$\frac{2EA}{L} \left(\begin{bmatrix} 2 & -1 \\ -1 & 1 \end{bmatrix} - \underbrace{\frac{\omega^2 \rho L^2}{24E}}_{\lambda} \begin{bmatrix} 4 & 1 \\ 1 & 2 \end{bmatrix} \right) \cdot \begin{Bmatrix} q_2 \\ q_3 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

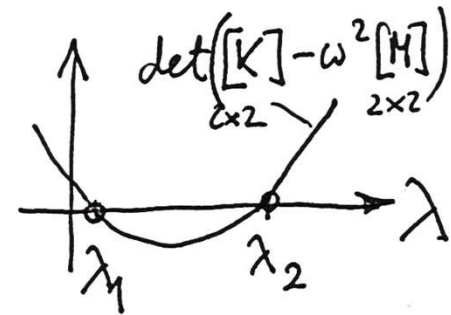
$$\left(\begin{bmatrix} 2 & -1 \\ -1 & 1 \end{bmatrix} - \lambda \begin{bmatrix} 4 & 1 \\ 1 & 2 \end{bmatrix} \right) \begin{Bmatrix} q_2 \\ q_3 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

$$\det \begin{bmatrix} 2-4\lambda & -(1+\lambda) \\ -(1+\lambda) & 1-2\lambda \end{bmatrix} = 0$$

$$(2-4\lambda)(1-2\lambda) - (-(1+\lambda)) \cdot (-(1+\lambda)) = 0$$

$$2-4\lambda-4\lambda+8\lambda^2-1-2\lambda-\lambda^2=0$$

$$7\lambda^2-10\lambda+1=0 ; \Delta = (-10)^2 - 4 \cdot 1 \cdot 7 = 72 \Rightarrow \sqrt{\Delta} = 6\sqrt{2}$$



$$\lambda_1 = \frac{-(-10) - 6\sqrt{2}}{14} = 0.1082, \quad \lambda_2 = \frac{-(-10) + 6\sqrt{2}}{14} = 1.3204$$

FIRST VIBRATION MODE :

$$\omega_1 = \sqrt{\frac{24E\lambda_1}{\rho L^2}} = \sqrt{\frac{24 \cdot 2 \cdot 10^5 \text{ MPa} \cdot 0.1082}{7.8 \cdot 10^{-9} \frac{\text{Ns}^2}{\text{mm}^4} \cdot (500\text{mm})^2}} = 16320 \frac{1}{\text{s}}$$

$$f_1 = \frac{\omega_1}{2\pi} = 2597.4 \text{ Hz}$$

Analytical solution: $\bar{\omega}_i = \frac{(i-0.5)\pi}{L} \sqrt{\frac{E}{\rho}}$, $\bar{\omega}_1 = \frac{\pi}{2L} \sqrt{\frac{E}{\rho}}$

Relative error: $\Delta\omega_1 = \frac{\omega_1 - \bar{\omega}_1}{\bar{\omega}_1} = 2.6\%$

SECOND VIBRATION MODE :

$$\omega_2 = \sqrt{\frac{24E\lambda_2}{\rho L^2}} = 57011 \frac{1}{\text{s}}, \quad f_2 = \frac{\omega_2}{2\pi} = 9073.5 \text{ Hz}$$

$$\bar{\omega}_2 = \frac{3\pi}{2L} \sqrt{\frac{E}{\rho}}, \quad \Delta\omega_2 = \frac{\omega_2 - \bar{\omega}_2}{\bar{\omega}_2} = 19.5\%$$

EIGENVECTORS: ($i=1,2$):

$$\left(\begin{bmatrix} 2 & -1 \\ -1 & 1 \end{bmatrix} - \lambda_i \begin{bmatrix} 4 & 1 \\ 1 & 2 \end{bmatrix} \right) \begin{Bmatrix} q_2(\lambda_i) \\ q_3(\lambda_i) \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

$$\begin{cases} (2 - 4\lambda_i) \cdot q_2(\lambda_i) - (1 + \lambda_i) \cdot q_3(\lambda_i) = 0 \\ -(1 + \lambda_i) \cdot q_2(\lambda_i) + (1 - 2\lambda_i) \cdot q_3(\lambda_i) = 0 \end{cases}$$

two linearly dependent equations

$$q_2(\lambda_i) = \frac{1 + \lambda_i}{2 - 4\lambda_i} \cdot q_3(\lambda_i)$$

NORMALIZATION

$$\underset{1 \times 3}{Lq} \cdot \underset{3 \times 3}{[I]} \cdot \underset{3 \times 1}{\{q\}} = 1$$

$$L0, q_2(\lambda_i), q_3(\lambda_i) \cdot \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \cdot \begin{Bmatrix} 0 \\ q_2(\lambda_i) \\ q_3(\lambda_i) \end{Bmatrix} = 1$$

$$L0, q_2(\lambda_i), q_3(\lambda_i) \cdot \begin{Bmatrix} 0 \\ q_2(\lambda_i) \\ q_3(\lambda_i) \end{Bmatrix} = 1$$

$$(q_2(\lambda_i))^2 + (q_3(\lambda_i))^2 = 1$$

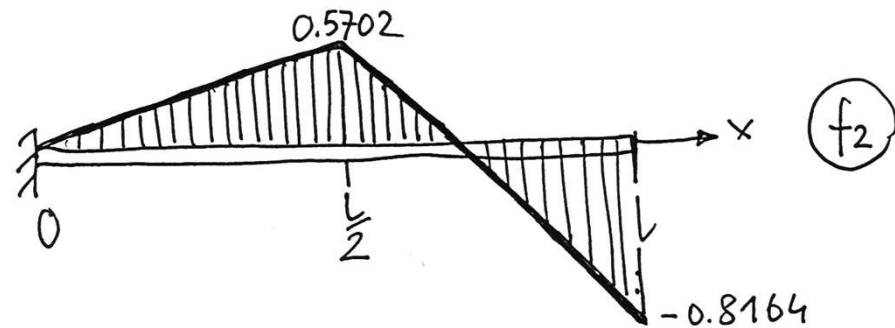
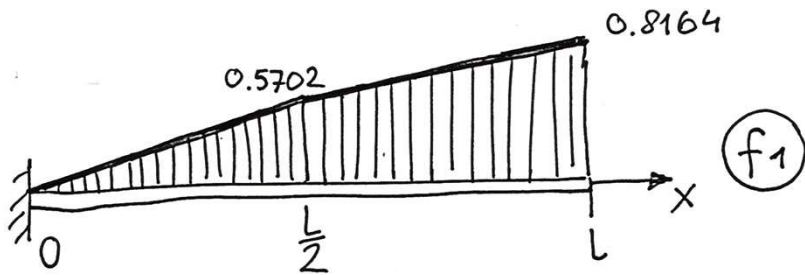
$$\frac{(1+\lambda_i)^2}{(2-4\lambda_i)^2} \cdot (q_3(\lambda_i))^2 + (q_3(\lambda_i))^2 = 1$$

$$\frac{(1+\lambda_i)^2 + (2-4\lambda_i)^2}{(2-4\lambda_i)^2} \cdot (q_3(\lambda_i))^2 = 1$$

$$q_3(\lambda_i) = \frac{2 - 4\lambda_i}{\sqrt{(1 + \lambda_i)^2 + (2 - 4\lambda_i)^2}}$$

$$q_2(\lambda_i) = \frac{1 + \lambda_i}{\sqrt{(1 + \lambda_i)^2 + (2 - 4\lambda_i)^2}}$$

i	λ_i	$q_2(\lambda_i)$	$q_3(\lambda_i)$	f_i [Hz]
1	0.1082	0.5702	0.8164	2597.4
2	1.3204	0.5702	-0.8164	9073.5



ANSYS:



	<i>1st mode</i>		<i>2nd mode</i>	
theory	2531.854 Hz	relative error	7595.545 Hz	relative error
2FEs	2597.294 Hz	2.6%	9074.278 Hz	19.5%
10FEs	2534.5 Hz	0.1%	7666 Hz	0.9%